## Aircraft Engines

Lecture Notes
(First Preliminary Edition)

## Section A. N omenclature Topics

## Part I. Thermodynamic Processes, Real and Ideal Gases

## Chapter 3. Theoretical Dependencies for an Ideal Gas

## Lecture 4. Parameters of an Ideal Gas in Interrelationships

## § 1.5. The equation of state for an ideal gas

An equation that establishes (sets up) the connection between pressure, temperature, and specific volume of a medium of a constant content (composition, compound) is called the thermodynamic equation of state [113, pp. 14-15], [79, p. 381]. The general view of this equation is [113, p. 14, (2.6)]

$$
\begin{equation*}
f(p, v, T)=0 . \tag{2.6}
\end{equation*}
$$

The equation of (2.6) in space depicts (portrays, shows, represents, describes, illustrates, gives a picture of) a surface, which characterises all possible equilibrium states of a regular (uniform) thermodynamic system. This surface is called the thermodynamic surface or the surface of state. At the thermodynamic surface a certain point corresponds to each state of the system.

In the equation of (2.6) the number of the independent variables will be equal to two of the three of $(p, v, T)$, because any of the three parameters (state variables) is the one-valued function of the other two given. For example (instance), if one accepts $v$ and $T$ as the independent variables, then $p$ can be expressed as a function of $v$ and $T$, i.e. (that is) $p=f_{p}(v, T)$ (the lower index (subscript) has a sense of the analytical view of the function of $f_{p}$ to get the $p$ as the dependent variable
(function); not to be confused with the constant state variables given below for isoprocesses and partial derivatives); if one accepts $p$ and $T$ as the independent variables, then specific volume $v=f_{v}(p, T)$.

If one of the parameters of state (state variables) is the constant value, then there will be two variable values, and points, representing the state of the system, will lay (lie) upon one plane, which intersects the thermodynamic surface perpendicularly to the coordinate axis, upon which the constant value is taken. Such coordinate systems on plane are called diagrams of state of substance. The diagrams with coordinates $p v, p T, v T, T s$, is are applied (used) most often.

The equation of state for the ideal gas first was obtained (found, got, deduced) by Benoît Paul Émile Clapeyron in 1834 by the way of joining (unifying, combining) the equations of the Robert's Boyle's-Edme's Mariotte's and Joseph's Louis's Gay-Lussac's laws

$$
\frac{p v}{T}=\text { const } .
$$

Denoting (indicate, designating, specifying) the constant as $R$, we will get [113, p. 14, (2.7)]

$$
\begin{equation*}
p v=R T \tag{2.7}
\end{equation*}
$$

where $R$ - specific gas constant, relatively to the mass of the gas is equaled to 1 kg ; $\mathrm{J} /(\mathrm{kg} \cdot \mathrm{K})$.

Benoît Paul Émile Clapeyron (26 February 1799 - 28 January 1864) was a French engineer and physicist, one of the founders of thermodynamics.

Robert Boyle, Fellow of the Royal Society (FRS), (25 January 1627 - 31 December 1691) was a $17^{\text {th }}$ century natural philosopher, chemist, physicist, and inventor, also noted for his writings in theology. He has been variously described as Irish, English and Anglo-Irish, his father having come to Ireland from England during the time of the Plantations.

Edme Mariotte (Til-Châtel c. 1620 - Paris 12 May 1684) was a French physicist and priest.

Joseph Louis Gay-Lussac (pronounced: [3วsعf lwi gelysak]; also Louis Joseph Gay-Lussac, 6 December 1778 - 9 May 1850) was a French chemist and physicist. He is known mostly for two laws related to gases, and for his work on alcohol-water mixtures, which led to the degrees Gay-Lussac used to measure alcoholic beverages in many countries.

One of the possible ways to deduce (derive, get, gain, obtain, receive, take, develop, originate) the equation of state (2.7) is the next (following). The Boyle'sMariotte's law (R. Boyle (1662), E. Mariotte (1676)) is

$$
p v=\text { const }, \quad T=\text { const ; }
$$

and the Gay-Lussac's law (1802) is

$$
\frac{V}{T}=\text { const }, \quad p=\mathrm{const} .
$$

I. Let us consider a chain of two of iso-processes that transfer the system from the state 1 with the coordinates of $\left(v_{1}, p_{1}\right)$, through the intermediate state $T p$ with the coordinates of $\left(v_{N_{2} 1}, p_{2}\right)$, to the state 2 with the coordinates of $\left(v_{2}, p_{2}\right)$ (fig. 2.a3). The iso-curves are plotted for nitrogen with the following (next) parameters (the specific gas constant see below, at the end of this paragraph). [XXX, pp. XXX-XXXX]

| Pa | $\frac{\mathrm{m}^{3}}{\mathrm{~kg}}$ | K |  |
| :---: | :---: | :---: | :---: |
| $\mathrm{p} 1:=1.01332 \cdot 10^{5}$ | $\mathrm{v} 1:=\frac{22.4146 \cdot 10^{-3}}{\mu \mathrm{~N} 2}$ | $\mathrm{v} 1=0.79995$ | $\mathrm{~T} 1:=\frac{\mathrm{p} 1 \cdot \mathrm{v} 1}{\mathrm{RN} 2}$ |
| $\mathrm{p} 2:=8 \cdot \mathrm{p} 1$ | $\mathrm{p} 2=8.107 \times 10^{5}$ | $\mathrm{v} 2:=\frac{\mathrm{v} 1}{2}$ | $\mathrm{v} 2=0.39998$ |
| $\mathrm{~T} 2:=\frac{\mathrm{p} 2 \cdot \mathrm{v} 2}{\mathrm{RN} 2}$ | $\mathrm{~T} 2=1.093 \times 10^{3}$ |  |  |

First, we compress the gas in the cylinder at the constant temperature $T_{1}$; carry out the compression and at the same time cooling the gas in order to keep the temperature of the gas constant (or doing the compression so slowly that the
temperature of the system relaxes). Then, at the constant pressure $p_{2}$ we expand the gas by heating it.


Fig. 2.a3

The first process $1-T p$ is the isothermal. Accordingly to the Boyle's-Mariotte's law

$$
\begin{equation*}
p v=\text { const }=p_{1} v_{1}=p_{2} v_{T p}, \quad T=\text { const }=T_{1} . \tag{2.7.1}
\end{equation*}
$$

The second process $T p-2$ is the isobaric. Accordingly to the Gay-Lussac's law

$$
\begin{equation*}
\frac{v}{T}=\text { const }=\frac{v_{T p}}{T_{1}}=\frac{v_{2}}{T_{2}}, \quad p=\text { const }=p_{2} \tag{2.7.2}
\end{equation*}
$$

Expressing $v_{T p}$ from the condition of (2.7.2) we get

$$
\begin{equation*}
v_{T p}=\frac{T_{1}}{T_{2}} v_{2} \tag{2.7.3}
\end{equation*}
$$

Substituting $v_{T p}$ of the equation of (2.7.3) into the relation of (2.7.1) we get

$$
\begin{equation*}
p_{1} v_{1}=p_{2} \frac{T_{1}}{T_{2}} v_{2}, \quad \frac{p_{1} v_{1}}{T_{1}}=\frac{p_{2} v_{2}}{T_{2}}=\text { const. } \tag{2.7.4}
\end{equation*}
$$

Thus, we come to the expression of (2.7), i.e. (that is) the equation of state by Clapeyron for the ideal gas.

Actually, the equation of (2.7) can be obtained by leading the process from the state 1 to the state 2 through the five other points.
II. For 1-Tv and $T v-2$

$$
\begin{equation*}
p v=\text { const }=p_{1} v_{1}=p_{T v} v_{2}, \quad T=\text { const }=T_{1} . \tag{2.7.5}
\end{equation*}
$$

The Charles's law (also the second Gay-Lussac's law (1808) is

$$
\begin{equation*}
\frac{p}{T}=\text { const }, \quad v=\text { const } . \tag{2.7.6}
\end{equation*}
$$

Then,

$$
\begin{gather*}
\frac{p}{T}=\text { const }=\frac{p_{T v}}{T_{1}}=\frac{p_{2}}{T_{2}}, \quad v=\text { const }=v_{2} .  \tag{2.7.7}\\
p_{T v}=p_{2} \frac{T_{1}}{T_{2}} .  \tag{2.7.8}\\
p_{1} v_{1}=p_{2} \frac{T_{1}}{T_{2}} v_{2}, \quad \frac{p_{1} v_{1}}{T_{1}}=\frac{p_{2} v_{2}}{T_{2}}=\text { const. } \tag{2.7.9}
\end{gather*}
$$

Jacques Alexandre César Charles (November 12, 1746 - April 7, 1823) was a French inventor, scientist, mathematician, and balloonist.
III. For $1-v p$ and $v p-2$

$$
\begin{array}{cc}
\frac{p}{T}=\text { const }=\frac{p_{1}}{T_{1}}=\frac{p_{2}}{T_{v p}}, \quad & v=\text { const }=v_{1} . \\
\frac{v}{T}=\text { const }=\frac{v_{1}}{T_{v p}}=\frac{v_{2}}{T_{2}}, \quad & p=\text { const }=p_{2} . \\
\frac{1}{T_{v p}}=\frac{v_{2}}{T_{2} v_{1}} . \\
\frac{p_{1}}{T_{1}}=\frac{p_{2} v_{2}}{T_{2} v_{1}}, & \frac{p_{1} v_{1}}{T_{1}}=\frac{p_{2} v_{2}}{T_{2}}=\text { const. } \tag{2.7.13}
\end{array}
$$

IV. For $1-v T$ and $v T-2$

$$
\begin{gather*}
\frac{p}{T}=\text { const }=\frac{p_{1}}{T_{1}}=\frac{p_{v T}}{T_{2}}, \quad v=\text { const }=v_{1} .  \tag{2.7.14}\\
p v=\text { const }=p_{v T} v_{1}=p_{2} v_{2}, \quad T=\text { const }=T_{2} .  \tag{2.7.15}\\
p_{v T}=\frac{p_{2} v_{2}}{v_{1}} .  \tag{2.7.16}\\
\frac{p_{1}}{T_{1}}=\frac{p_{2} v_{2}}{T_{2} v_{1}}, \quad \frac{p_{1} v_{1}}{T_{1}}=\frac{p_{2} v_{2}}{T_{2}}=\text { const. } \tag{2.7.17}
\end{gather*}
$$

V. For $1-p v$ and $p v-2$

$$
\begin{gather*}
\frac{v}{T}=\text { const }=\frac{v_{1}}{T_{1}}=\frac{v_{2}}{T_{p v}}, \quad p=\text { const }=p_{1} .  \tag{2.7.18}\\
\frac{p}{T}=\text { const }=\frac{p_{1}}{T_{p v}}=\frac{p_{2}}{T_{2}}, \quad v=\text { const }=v_{2} .  \tag{2.7.19}\\
\frac{1}{T_{p v}}=\frac{p_{2}}{T_{2} p_{1}} .  \tag{2.7.20}\\
\frac{v_{1}}{T_{1}}=\frac{v_{2} p_{2}}{T_{2} p_{1}}, \quad \frac{v_{1} p_{1}}{T_{1}}=\frac{v_{2} p_{2}}{T_{2}}=\text { const } . \tag{2.7.21}
\end{gather*}
$$

VI. For $1-p T$ and $p T-2$

$$
\begin{gather*}
\frac{v}{T}=\text { const }=\frac{v_{1}}{T_{1}}=\frac{v_{p T}}{T_{2}}, \quad p=\text { const }=p_{1} .  \tag{2.7.22}\\
p v=\text { const }=p_{1} v_{p T}=p_{2} v_{2}, \quad T=\text { const }=T_{2} .  \tag{2.7.23}\\
v_{p T}=\frac{p_{2} v_{2}}{p_{1}} .  \tag{2.7.24}\\
\frac{v_{1}}{T_{1}}=\frac{v_{2} p_{2}}{T_{2} p_{1}}, \quad \frac{v_{1} p_{1}}{T_{1}}=\frac{v_{2} p_{2}}{T_{2}}=\text { const. } \tag{2.7.25}
\end{gather*}
$$

The equation of (2.7) is put (written) down for 1 kg of an ideal gas. For $m \mathrm{~kg}$, the equation of state will have the view of [113, p. 15, (2.8)]

$$
\begin{equation*}
p V=m R T \tag{2.8}
\end{equation*}
$$

where $V$ - the volume of the gas, $\mathrm{m}^{3}$.
The gas which state is accurately described by the equation of (2.7) is called the ideal.

Multiplying the both parts of the equation of (2.7) by the molecular mass (weight) $\mu$, we will get [113, p. 15, (2.9)]

$$
\begin{equation*}
p V_{\mu}=\mu R T \tag{2.9}
\end{equation*}
$$

where $V_{\mu}=v \mu$ - the volume occupied by one mole of the gas.
Molecular mass (weight) represents

The volume of one mole of gas at the normal physical conditions $\left(T_{0}=273.15 \mathrm{~K}, p_{0}=101,332 \mathrm{~Pa}\right) V_{\mu}=\mu \nu=22.4146 \cdot 10^{-3} \mathrm{~m}^{3} / \mathrm{mole}$. Substituting this value into the formula of (2.9) we will get [113, p. 15, (2.9a)]

$$
\begin{equation*}
\mu R=\frac{p_{0} V_{\mu}}{T_{0}}=8.3143 \mathrm{~J} /(\mathrm{mole} \cdot \mathrm{~K}) \tag{2.9a}
\end{equation*}
$$

where $\mu R$ - universal gas constant, the same for any of gases [79, p. 381].
From here the equation of state for one mole of the idea gas will be [113, p. 15, (2.10)]

$$
\begin{equation*}
p V_{\mu}=8.3143 T \tag{2.10}
\end{equation*}
$$

The equation of (2.10) was deduced by D.I. Mendeleyev in 1874 and is called the equation of state by Clapeyron-Mendeleyev [113, p. 14-15].

The molecular masses (weights) of hydrogen, nitrogen, and oxygen [XXX, pp. XXX-XXXX]

$$
\mu \mathrm{H} 2:=2.016 \cdot 10^{-3} \quad \mu \mathrm{~N} 2:=28.02 \cdot 10^{-3} \quad \mu \mathrm{O} 2:=32 \cdot 10^{-3} \quad \frac{\mathrm{~kg}}{\mathrm{~mol}}
$$

The specific gas constants for hydrogen, nitrogen, and oxygen

$$
\begin{array}{llll}
\mathrm{RH} 2:=\frac{8.3143}{\mu \mathrm{H} 2} & \mathrm{RN} 2:=\frac{8.3143}{\mu \mathrm{~N} 2} & \mathrm{RO} 2:=\frac{8.3143}{\mu \mathrm{O} 2} & \frac{\frac{\mathrm{Joules}}{\mathrm{molK}}}{\frac{\mathrm{~kg}}{\mathrm{~mol}}} \\
\mathrm{RH} 2=4124.157 & \mathrm{RN} 2=296.727 & \mathrm{RO} 2=259.822 & \frac{\mathrm{Joules}}{\mathrm{~kg} \cdot \mathrm{~K}}
\end{array}
$$

